

19 April 2021

THERE IS ONLY ONE TIME

time and classical equations of motion from
quantum entanglement via the
Page and Wootters mechanism
with generalized coherent states

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unifi & infn

SUMMARY

★ WHAT COMES FIRST

★ THIS WORK

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- i1. is there a problem ?
- i2. GCS and the large- N limit
- i3. the PAW mechanism

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★ THIS WORK

- H1. a quantum clock for a quantum system QQ
→ energy-time uncertainty relation
- H2. a classical clock for a quantum system CQ
→ Schrödinger equation
- H3. a classical clock for a classical system CC
→ Hamilton equations of motion !

IS THERE A PROBLEM ?

- why should there be any
spacetime is cool
(and QFT is amazing)

IS THERE A PROBLEM ?

- why should there be any
spacetime is cool
(and QFT is amazing)
- ... uhm ...
- time is not a quantum observable
 $\Delta E \Delta t \geq \hbar$ does not make sense
(and QFT does not help when it comes to

GR \longleftrightarrow QM)

?

GENERALIZED COHERENT STATES (GCS)

$\mathcal{H}, \mathfrak{g}, |q\rangle$

$|\alpha\rangle \in \mathcal{H} \leftrightarrow \alpha \in \mathcal{H}$

GENERALIZED COHERENT STATES (GCS)

$\mathcal{H}, \mathfrak{g}, |q\rangle$

$$|\alpha\rangle \in \mathcal{H} \leftrightarrow \alpha \in \mathcal{M}$$

$$\int_{\mathcal{M}} \phi(\hat{\alpha}) |\alpha \times \alpha| = \hat{\mathbb{1}}_{\mathcal{H}}, \quad d\mu(\hat{\alpha}) \text{ invariant}, \quad \langle \alpha | \alpha' \rangle \neq \delta(\alpha - \alpha')$$

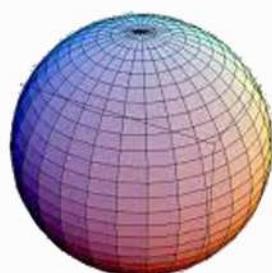
GENERALIZED COHERENT STATES (GCS)

$\mathcal{H}, \mathfrak{g}, |\zeta\rangle$

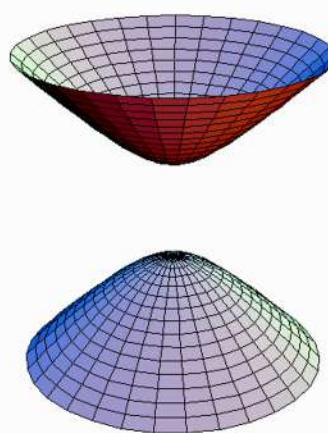
$|\zeta\rangle \in \mathcal{H} \leftrightarrow \zeta \in \mathcal{M}$

$\int_{\mathcal{M}} \phi(\hat{\mu}) |\zeta \times \zeta| = \hat{L}_{\mathcal{H}}, \phi(\hat{\mu})$ invariant, $\langle \zeta | \zeta' \rangle \neq \delta(\zeta - \zeta')$

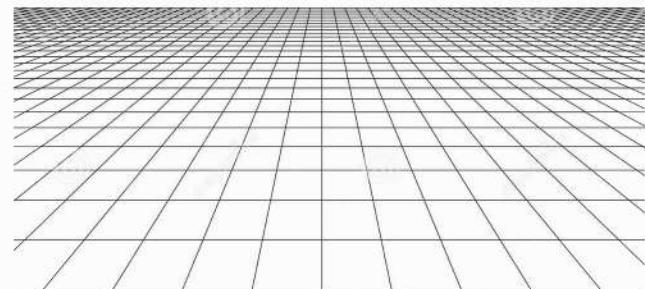
\mathcal{M} is symplectic and depends on \mathfrak{g}



$\mathfrak{su}(2)$



$\mathfrak{su}(1,1)$



\mathbb{h}_4

GCS AND THE LARGE-N LIMIT

QUANTUM-TO-CLASSICAL CROSSOVER

conditions : $Q_N \xrightarrow[N \rightarrow \infty]{} \alpha$, in terms of GCS

GCS AND THE LARGE-N LIMIT

QUANTUM-TO-CLASSICAL CROSSOVER

conditions: $Q_N \xrightarrow[N \rightarrow \infty]{} \Omega$, in terms of GCS



states $|\psi\rangle$ that survive:

$$\langle \omega | \omega' \rangle \xrightarrow[N \rightarrow \infty]{} \delta(\omega - \omega')$$

$$|\omega\rangle \in \mathcal{H}$$

$$\omega \in \mathcal{M}$$

observables \hat{A} that survive:

$$\frac{\langle \omega | \hat{A} | \omega' \rangle}{\langle \omega | \omega' \rangle} \xrightarrow[N \rightarrow \infty]{} \infty$$

$$\langle \omega | \hat{A} | \omega \rangle$$

$$A(\omega)$$

symbol

GCS AND COMPOSITE SYSTEMS $\Psi = C + F$

A PARAMETRIC REPRESENTATION

$$|\Psi\rangle = \sum_{\gamma} c_{\gamma} |\gamma\rangle \otimes |y\rangle \in \mathcal{H}_{\Psi} = \mathcal{H}_c \otimes \mathcal{H}_F$$

GCS AND COMPOSITE SYSTEMS $\Psi = C + F$

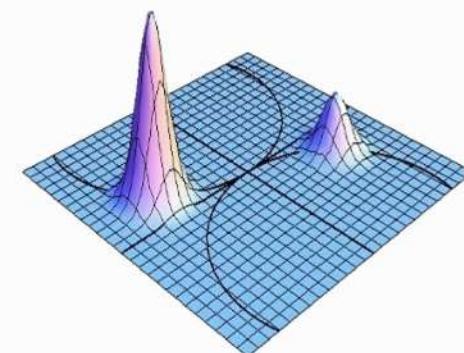
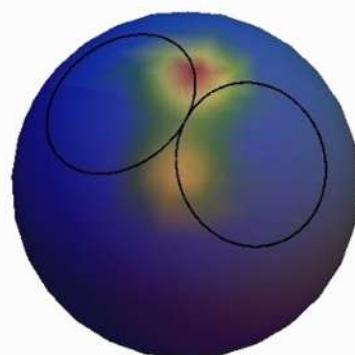
A PARAMETRIC REPRESENTATION

$$\begin{aligned}
 |\Psi\rangle = & \sum_{\lambda} c_{\lambda} |\lambda\rangle \otimes |\gamma\rangle \in \mathcal{H}_{\Psi} = \mathcal{H}_e \otimes \mathcal{H}_F \\
 & \underbrace{\qquad\qquad\qquad}_{\int d\mu(\omega) |e\rangle |\chi(\omega)\rangle} \\
 = & \int_{\mathcal{M}} d\mu(\omega) |\chi(\omega)\rangle \otimes |\phi(\omega)\rangle
 \end{aligned}$$

GCS AND COMPOSITE SYSTEMS $\Psi = C + F$

A PARAMETRIC REPRESENTATION

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 & \underbrace{\qquad\qquad\qquad}_{\int d\mu(\omega) |e\rangle \chi(\omega)} \\
 = & \int_{\mathcal{M}} d\mu(\omega) |\chi(\omega)\rangle |\Omega\rangle \otimes |\phi(\omega)\rangle \\
 & \chi^2(\omega) = \sum_{\Omega} \left| \sum_{\chi} c_{\Omega\chi} \langle \Omega | \chi \rangle \right|^2
 \end{aligned}$$



THE PAW) MECHANISM

@ a certain time t :

conditioned to a clock being in a state labeled by t

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conditioned to a clock being in a state labeled by t



in a quantum setting: $\Psi = C + \Gamma$ and 3 assumptions

1.

$$\hat{H} = \hat{H}_c \otimes \hat{1}_r - \hat{1}_c \otimes \hat{H}_r \quad \text{non interacting}$$

2.

$$|\Psi\rangle = \int_{\Omega} d\mu(\omega) \chi^2(\omega) |\omega\rangle \otimes |\phi(\omega)\rangle \quad \text{entangled}$$

3.

$$\hat{H} |\Psi\rangle = E |\Psi\rangle \quad E = 0 \text{ if you like}$$

GENERALIZED CS + RAW MECHANISM

algebras

hamiltonians

GENERALIZED CS + RAW MECHANISM

algebras

hamiltonians

Cartan decomposition

diagonal & shift

$$\hat{D}_s$$

$$\hat{R}_m$$

$$\hat{R}_m^+ = \hat{R}_{-m}$$

$$[\hat{D}_s, \hat{D}_\theta] = 0 \quad ; \quad [\hat{D}_s, \hat{R}_m] = d_{sm} \hat{R}_m \quad ; \quad [\hat{R}_m, \hat{R}_{-m}] = \sum_s d_{sm} \hat{D}_s \quad ; \quad [\hat{R}_m, \hat{R}_{m'}] = c_{mm'} \hat{R}_{m+m'}$$

GENERALIZED CS + RAW MECHANISM

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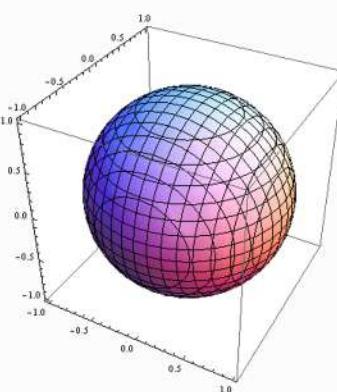
$$\hat{D}_\delta$$

$$\hat{R}_m$$

$$\hat{R}_m = \hat{R}_{-m}$$

$$[\hat{D}_\delta, \hat{D}_\theta] = 0 \quad ; \quad [\hat{D}_\delta, \hat{R}_m] = d_{\delta m} \hat{R}_m \quad ; \quad [\hat{R}_m, \hat{R}_{-m}] = \sum_\delta d_{\delta m} \hat{D}_\delta \quad ; \quad [\hat{R}_m, \hat{R}_{m'}] = c_{mm'} \hat{R}_{m+m'}$$

keep it simple : $\delta = m = 1$ and $d_{1, \pm 1} = \epsilon \in \mathbb{R}$ as in $\text{su}(2)$



reference state $|G\rangle : \hat{R}|G\rangle = 0 \Rightarrow \hat{D}|G\rangle \sim |G\rangle$

GENERALIZED CS + RAW MECHANISM

algebras



hamiltonians

- Cartan decomposition
diagonal & shift

$$[\hat{D}, \hat{R}] = \epsilon \hat{R} \quad [\hat{R}, \hat{R}^+] = \epsilon \hat{D}$$

- reference state $\hat{R}|G\rangle = 0$

$$|\lambda\rangle = e^{\lambda\hat{R}^+ - \lambda^*\hat{R}} |G\rangle = N_\rho e^{\lambda\hat{R}^+} |G\rangle$$

$$\lambda = \rho e^{-i\varphi} \in \mathbb{C}$$

$$\Lambda = |\tan \rho| e^{-i\varphi}$$

GENERALIZED CS + RAW MECHANISM

algebras



hamiltonians

- Cartan decomposition
diagonal & shift

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take

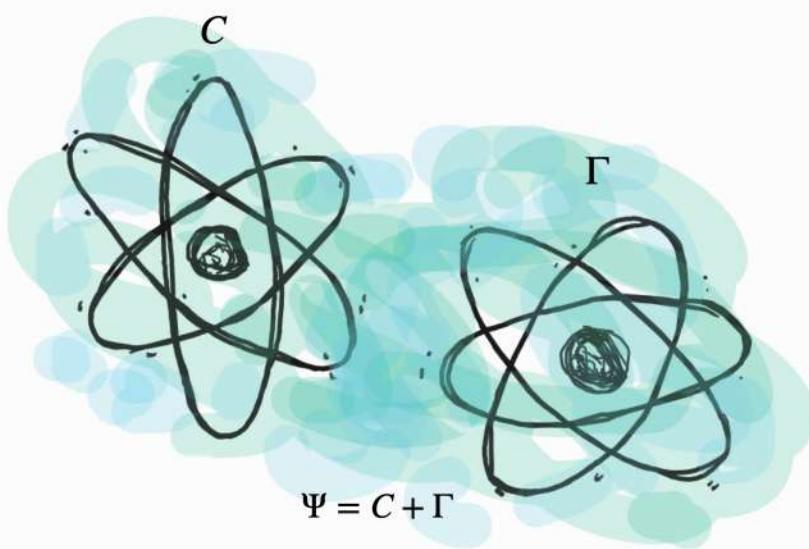
$$\hat{D} = \hat{H}_c$$

$$[\hat{H}_c, \hat{R}] = \epsilon \hat{R}$$

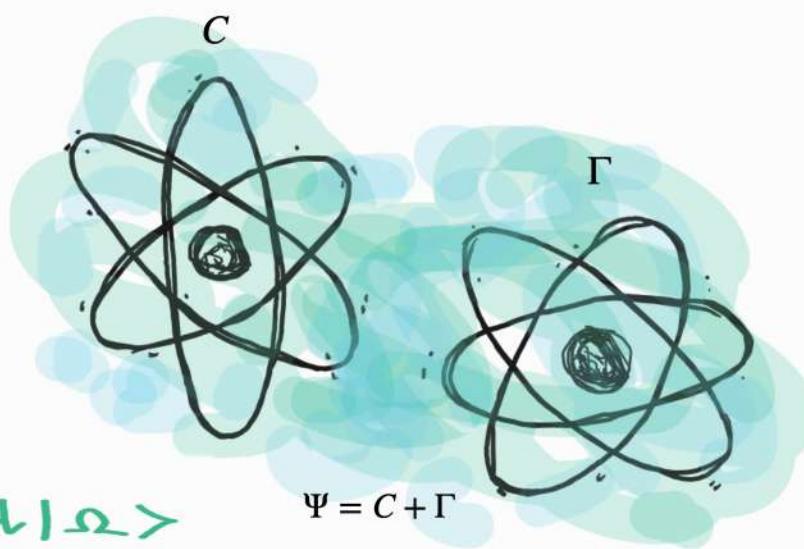
$$\hat{H}_c |G\rangle = \epsilon_0 |G\rangle = 0 \text{ if you like}$$

$$\hat{H}_c$$

A QUANTUM CLOCK FOR A QUANTUM SYSTEM



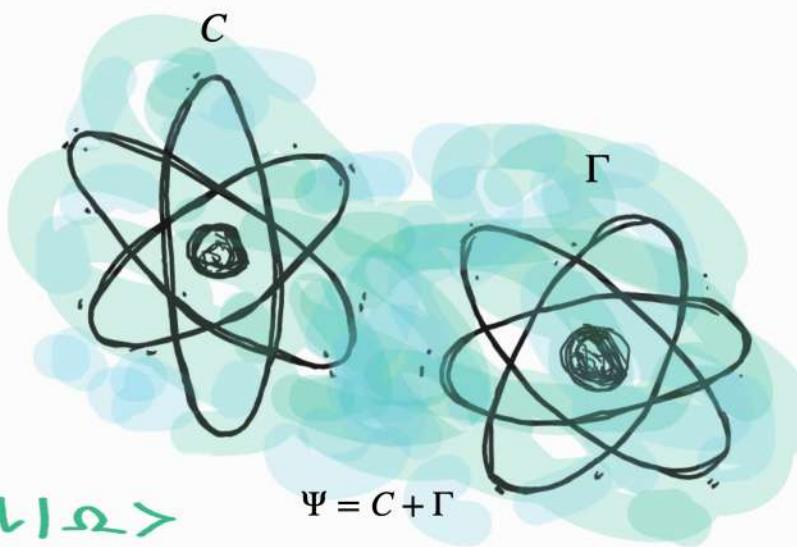
A QUANTUM CLOCK FOR A QUANTUM SYSTEM



$$[\hat{H}_c, e^{\lambda^* \hat{R}}] = \epsilon \lambda^* \hat{R} e^{\lambda^* \hat{R}}$$

$$\langle \lambda | \hat{H}_c | \Omega \rangle = i\epsilon \frac{d}{d\phi} \langle \lambda | \Omega \rangle$$

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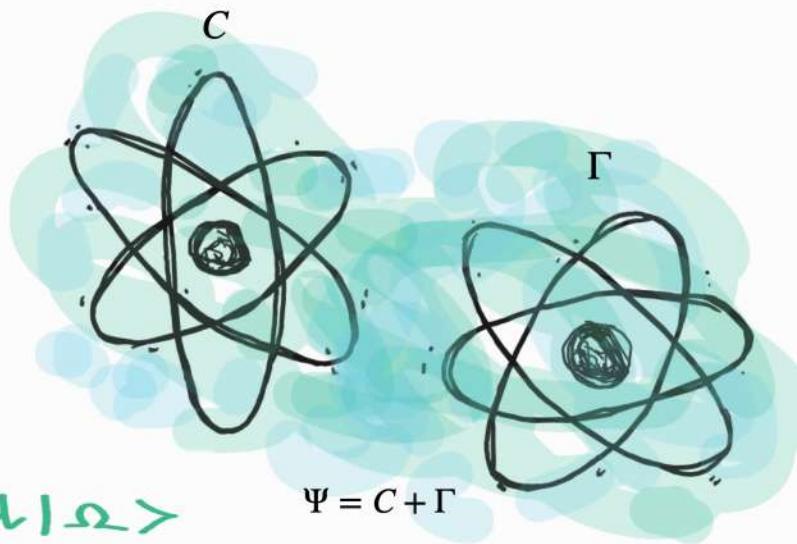
+

$$\langle \lambda | \hat{H} | \Psi \rangle = 0$$

↓

$$i\epsilon \frac{d}{d\varphi} |\Phi_\epsilon(\varphi)\rangle = \hat{H}_r |\Phi_\epsilon(\varphi)\rangle$$

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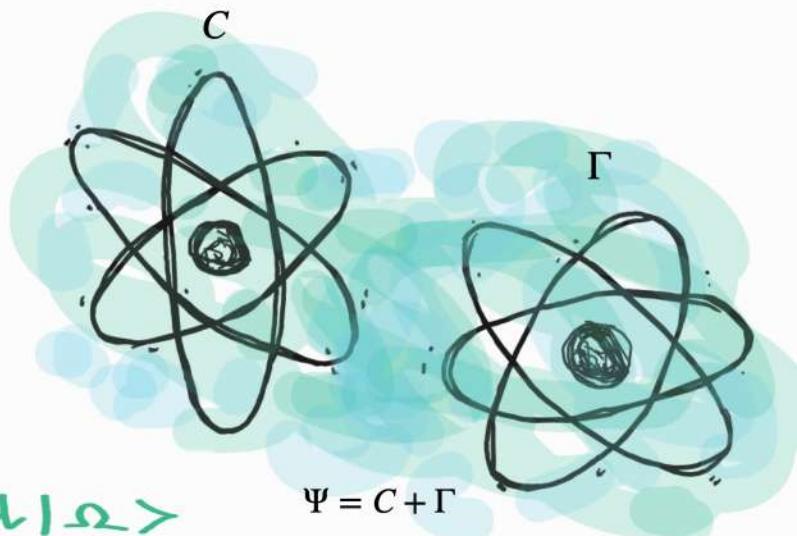
Schroed - like

$$\text{but } |\Phi_\epsilon(\varphi)\rangle = \langle \lambda | \Psi \rangle$$

unnormalized



A QUANTUM CLOCK FOR A QUANTUM SYSTEM



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$$\hat{R} = (\hat{R} \hat{R}^\dagger)^{1/2} e^{-i\hat{\phi}}$$

$$[\hat{H}_c, \sin \hat{\phi}] = i\epsilon \cos \hat{\phi}$$

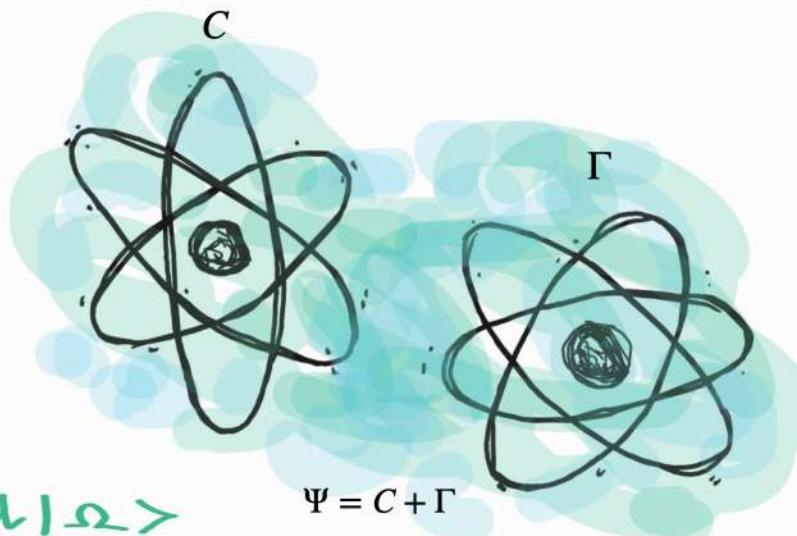
+

$$\hat{H}_c = \hat{H} + \hat{H}_r$$

↓

$$\Delta(\hat{H} + \hat{H}_r) \cdot \Delta \sin \hat{\phi} \geq \left| \frac{\epsilon}{2} \langle \cos \hat{\phi} \rangle \right|$$

A QUANTUM CLOCK FOR A QUANTUM SYSTEM



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$$[\hat{H}_c, \sin \hat{\phi}] = i\varepsilon \cos \hat{\phi}$$

$$[\hat{H}_c, e^{i\lambda^* \hat{R}}] = \varepsilon \lambda^* \hat{R} e^{i\lambda^* \hat{R}}$$

$$\langle \lambda | \hat{H}_c | \Omega \rangle = i\varepsilon \frac{d}{d\phi} \langle \lambda | \Omega \rangle$$

+

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?

+

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↓

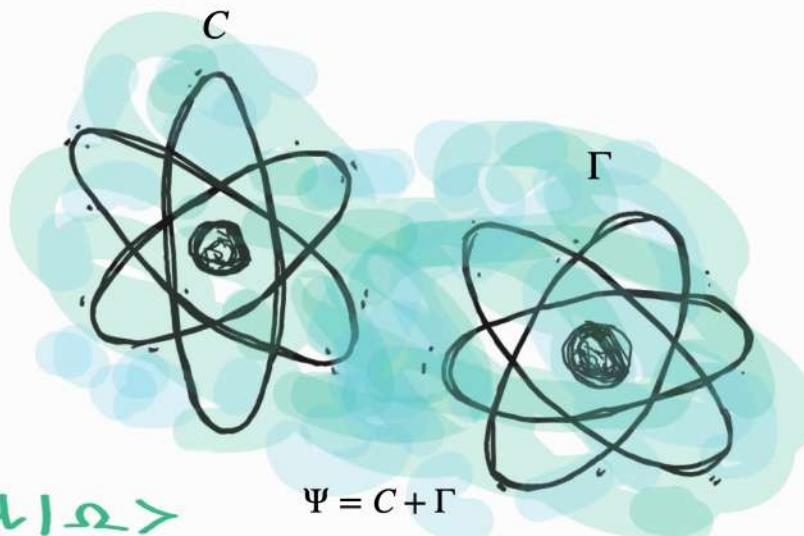
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energy-time-like

but not yet

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A QUANTUM CLOCK FOR A QUANTUM SYSTEM



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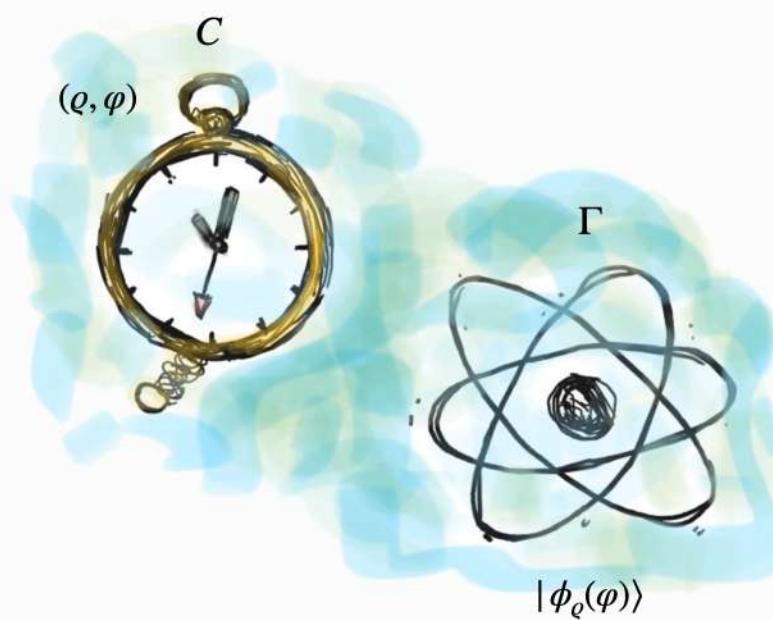
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energy-time-like
but not yet

$$\frac{\hbar}{\varepsilon} \varphi$$

A CLASSICAL CLOCK FOR A QUANTUM SYSTEM



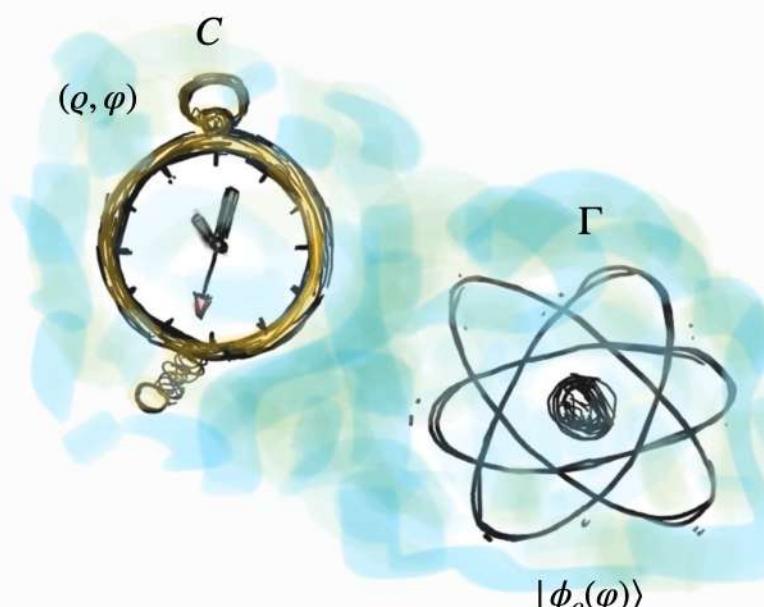
A CLASSICAL CLOCK FOR A QUANTUM SYSTEM

$$\lim_{N \rightarrow \infty} \langle \lambda | \omega \rangle = \delta(\lambda - \omega)$$

$$\langle \Phi_e(\varphi) | \Phi_e(\varphi) \rangle = \chi^2(e)$$

$$|\phi_e(\varphi)\rangle = \frac{|\Phi_e(\varphi)\rangle}{\chi(e)}$$

normalized



$$\text{i.e. } \frac{d}{d\varphi} |\phi_e(\varphi)\rangle = \hat{H}_\Gamma |\phi_e(\varphi)\rangle$$

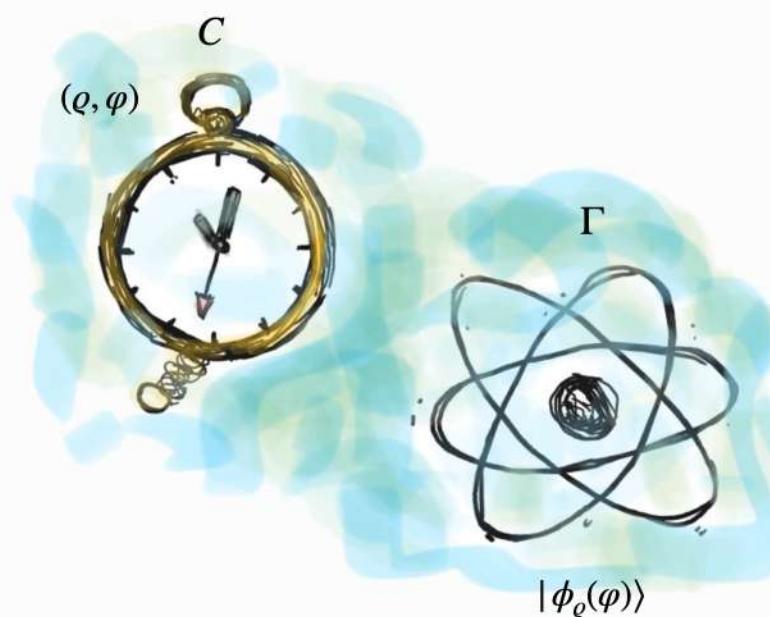
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$$\lim_{N \rightarrow \infty} \frac{\langle \lambda | \hat{A} | \omega \rangle}{\langle \lambda | \omega \rangle} < \infty$$

$$\langle \lambda | \hat{A} | \lambda \rangle = A(\lambda)$$

using

$$\langle \lambda | \hat{H} | \psi \rangle = 0$$

from PaW

$$\text{i.e. } \frac{d}{d\varphi} |\phi_e(\varphi)\rangle = \hat{H}_r |\phi_e(\varphi)\rangle$$

$$\hat{H}_r |\phi_e(\varphi)\rangle = E_r(e) |\phi_e(\varphi)\rangle$$

$$\begin{aligned} E_r(e) &= \langle \lambda | \hat{H}_c | \lambda \rangle \\ &= \frac{ek}{2} (1 - \cos 2e) \end{aligned}$$

$$\Delta E_r(e) \Delta \varphi \geq \frac{\varepsilon}{2}$$

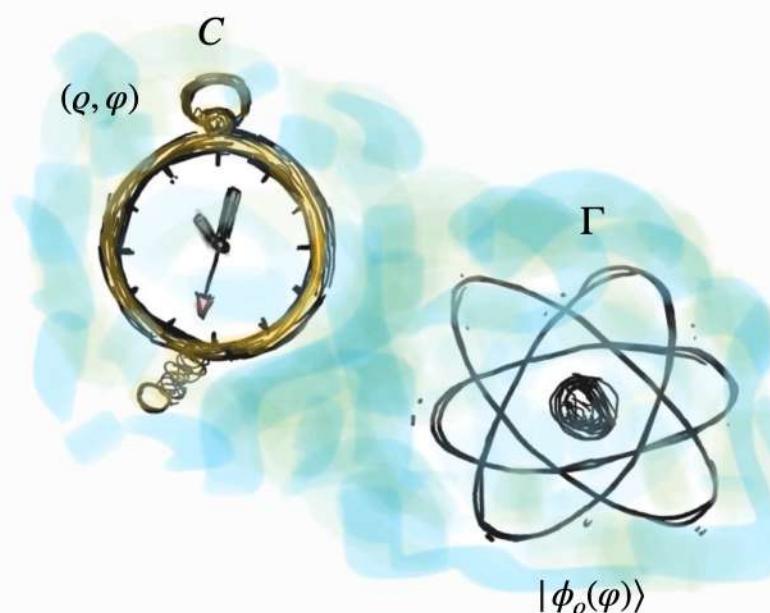
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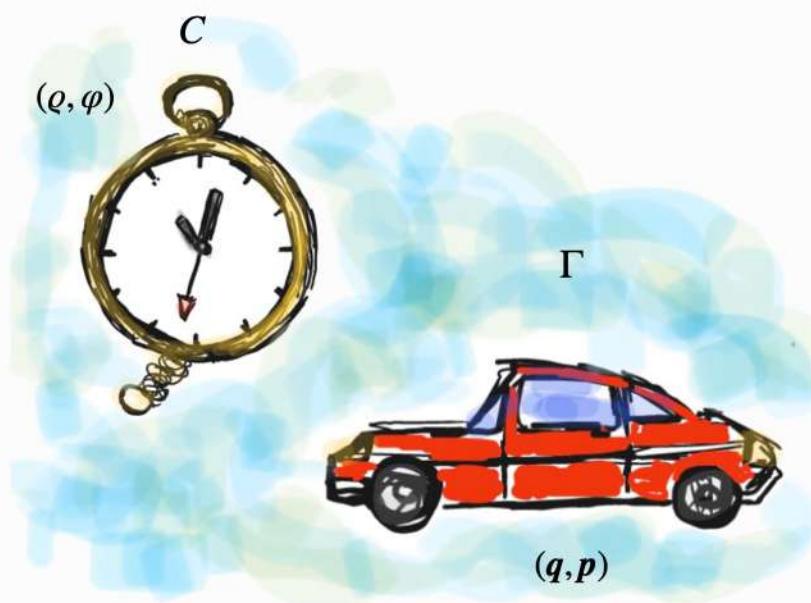
$$\begin{aligned} E_r(e) &= \langle \lambda | \hat{H}_c | \lambda \rangle \\ &= \frac{\epsilon k}{2} (1 - \cos 2e) \end{aligned}$$

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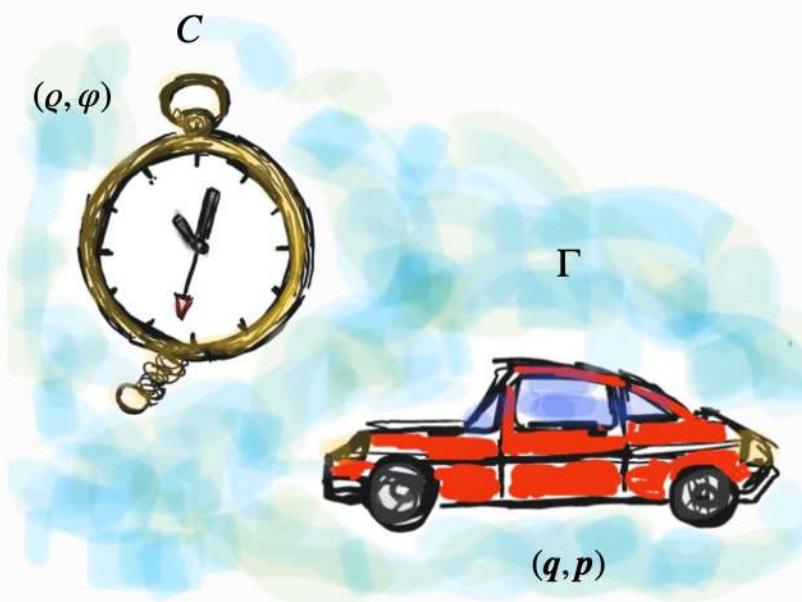
$$\dot{x}^{QM} = \frac{k}{\epsilon} \varphi$$

A CLASSICAL CLOCK FOR A CLASSICAL SYSTEM



A CLASSICAL CLOCK FOR A CLASSICAL SYSTEM

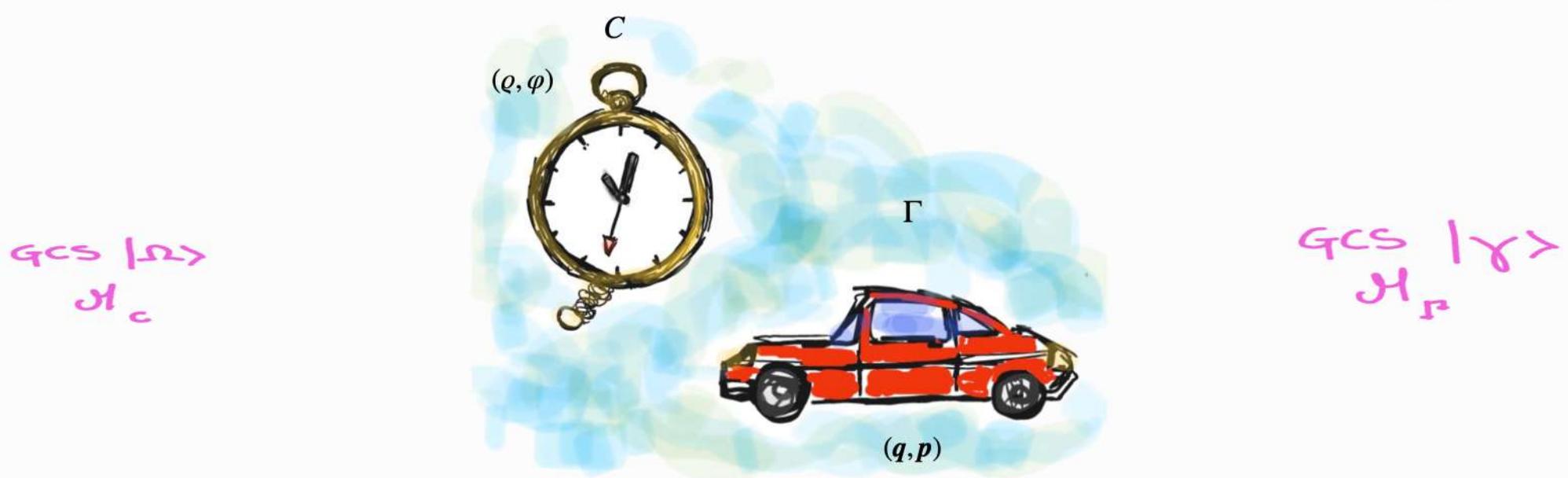
GCS $|\alpha\rangle$
 \mathcal{H}_c



GCS $|\gamma\rangle$
 \mathcal{H}_Γ

$$|\Psi\rangle = \int_{\mathcal{H}_c} d\mu(\alpha) \int_{\mathcal{H}_\Gamma} d\mu(\gamma) \Phi(\alpha, \gamma) |\alpha\rangle \otimes |\gamma\rangle$$

A CLASSICAL CLOCK FOR A CLASSICAL SYSTEM



$$|\Psi\rangle = \int_{\mathcal{H}_c} d\mu(\omega) \int_{\mathcal{H}_r} d\mu(\gamma) \beta(\omega, \gamma) |\omega\rangle \otimes |\gamma\rangle$$

+

using $\hat{H}|\Psi\rangle = 0$ from PaW



$$\langle \omega | H_c | \omega \rangle = \langle \gamma | H_r | \gamma \rangle \quad \text{for } \omega, \gamma : \beta(\omega, \gamma) \neq 0$$

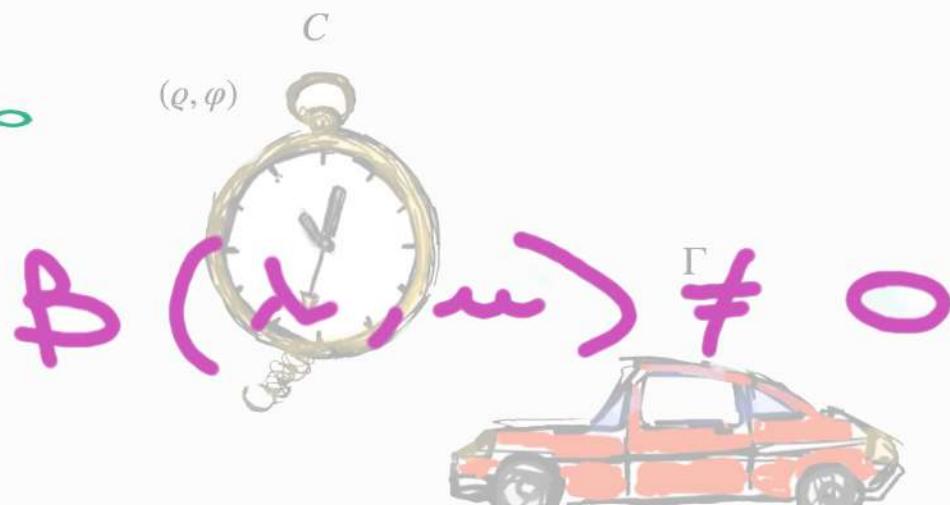
A CLASSICAL CLOCK FOR A CLASSICAL SYSTEM

$$\omega, \gamma: \beta(\omega, \gamma) \neq 0$$

$$\lambda, \mu$$

 \hookrightarrow

$$\lambda = \rho e^{-i\varphi}$$



$$H_c(\eta) = H_{\Gamma}(\mu)^{(q,p)}$$

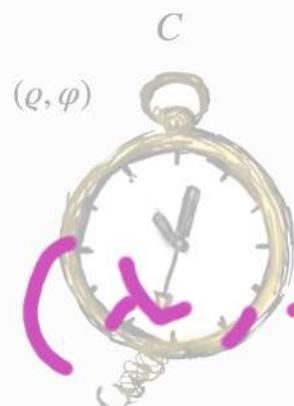
A CLASSICAL CLOCK FOR A CLASSICAL SYSTEM

$$\omega, \gamma: \beta(\omega, \gamma) \neq 0$$

$$\lambda, u$$

 \hookrightarrow

$$\lambda = \rho e^{-i\varphi}$$



$$\beta(\lambda, u) \neq 0$$



$$H_c(\rho) = H_r(u)$$



F:
:

$$H_c(\rho) = H_r(u = F(\lambda))$$

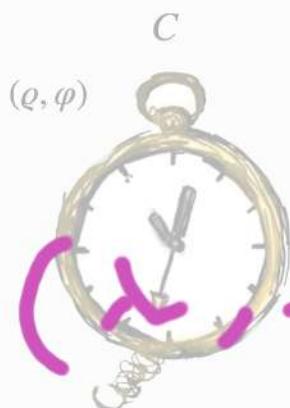
A CLASSICAL CLOCK FOR A CLASSICAL SYSTEM

$$\omega, \gamma: \beta(\omega, \gamma) \neq 0$$

$$\lambda, u$$

\$\hookrightarrow\$

$$\lambda = \rho e^{-i\varphi}$$



$$\beta(\lambda, u) \stackrel{\Gamma}{\neq} 0$$



$$H_c(\rho) = H_{\Gamma}(u)$$

$$\downarrow$$

$$F:$$

?

$$H_c(\rho) = H_{\Gamma}(u = F(\lambda))$$

H_{Γ} is symplectic

$\exists \beta: \gamma \rightarrow q, p /$

$$\{q, p\}_{\Gamma} = \frac{1}{\hbar}$$

for this pair
choose eq

$$q - ip = \sqrt{2\epsilon} \sin \varphi e^{-i\varphi}$$



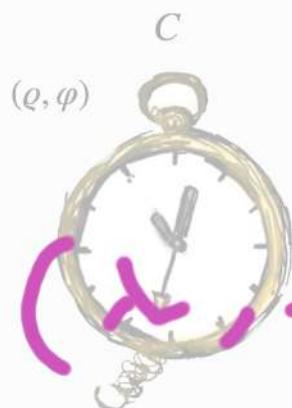
A CLASSICAL CLOCK FOR A CLASSICAL SYSTEM

$$\omega, \gamma: \beta(\omega, \gamma) \neq 0$$

$$\lambda, u$$

\$\hookrightarrow\$

$$\lambda = \rho e^{-i\varphi}$$



$$\beta(\lambda, u) \neq 0$$

$$H_c(\rho) = H_{\mu}(u)$$

$$\downarrow \\ F:$$

?

$$\frac{k\varepsilon}{2} (\cos 2\rho - 1) =$$

$$H_c(\rho) = H_{\mu}(u = F(\lambda))$$

$$= k\varepsilon(q^2 + p^2)$$

H_{μ} is symplectic

$\exists \beta: \gamma \rightarrow q, p /$

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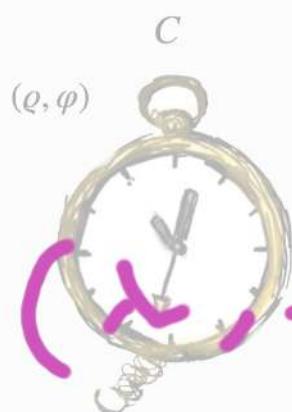
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$$\downarrow$$

F

?

$$\frac{k\varepsilon}{2} (\cos 2\rho - 1) =$$

$$H_c(\rho) = H_{\mu}(u = F(\lambda))$$

$$= k\varepsilon(q^2 + p^2)$$

pullback - by - F

$$\downarrow$$

F^*

H_{μ} is symplectic

$$\exists \beta: \gamma \rightarrow q, p /$$

$$\{q, p\}_{\mu} = \frac{1}{\hbar}$$

for this pair
choose eq

$$q - ip = \sqrt{2\varepsilon} \sin \rho e^{-i\varphi}$$

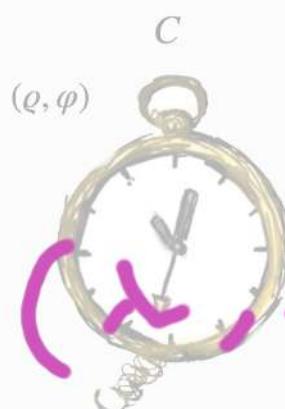


A CLASSICAL CLOCK FOR A CLASSICAL SYSTEM

$$\omega, \gamma: \beta(\omega, \gamma) \neq 0$$

$$\lambda, u \\ \hookrightarrow$$

$$\lambda = \rho e^{-i\varphi}$$



$$\beta(\lambda, u) \neq 0$$



$$H_c(\rho) = H_r(u)$$

$$\downarrow \\ F:$$

?

$$\frac{k\varepsilon}{2} (\cos 2\rho - 1) =$$

$$H_c(\rho) = H_r(u = F(\lambda))$$

$$= k\varepsilon(q^2 + p^2)$$

pullback - by - F

$$\downarrow \\ F^*$$

$$\{q, p\}_r \xrightarrow{F^*} \{f \cdot g\}_c$$

$$= (k\varepsilon \sin(2\rho))'$$

$$\frac{\partial f}{\partial \rho} \frac{\partial g}{\partial \varphi} - \frac{\partial g}{\partial \rho} \frac{\partial f}{\partial \varphi}$$

write

$$\{q, H_c\}_c \xrightarrow{F} \{q, H_r\}_r$$

and finally find

H_r is symplectic

$\exists \beta: \gamma \rightarrow q, p /$

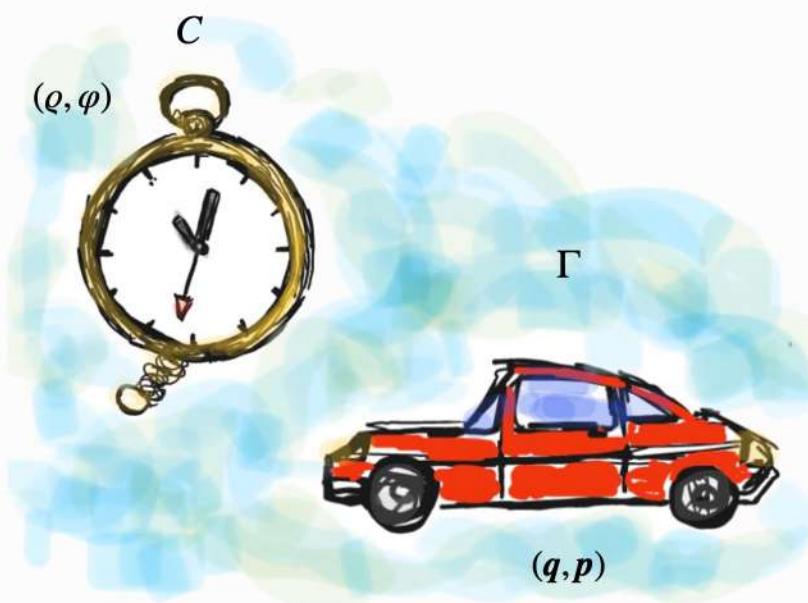
$$\{q, p\}_r = \frac{1}{\hbar}$$

for this pair
choose eq

$$q \cdot ip = \sqrt{2\varepsilon} \sin \rho e^{-i\varphi}$$



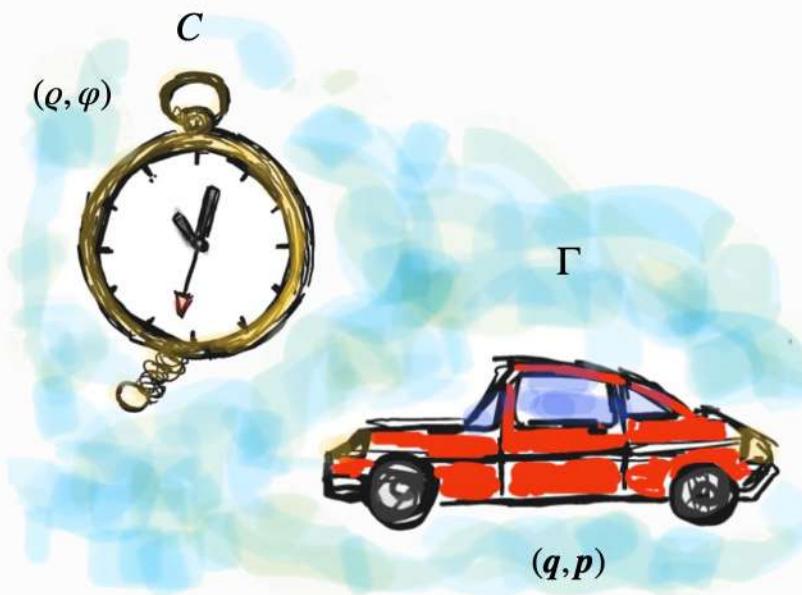
A CLASSICAL CLOCK FOR A CLASSICAL SYSTEM



$$\{ q, H_{\Gamma} \}_{\mu} = \frac{\epsilon}{\hbar} \frac{dq}{d\varphi}$$

$$\{ p, H_{\Gamma} \}_{\mu} = \frac{\epsilon}{\hbar} \frac{dp}{d\varphi}$$

A CLASSICAL CLOCK FOR A CLASSICAL SYSTEM



$$\{ q, H_{\Gamma} \}_p = \frac{\epsilon}{\hbar} \frac{dq}{d\varphi}$$

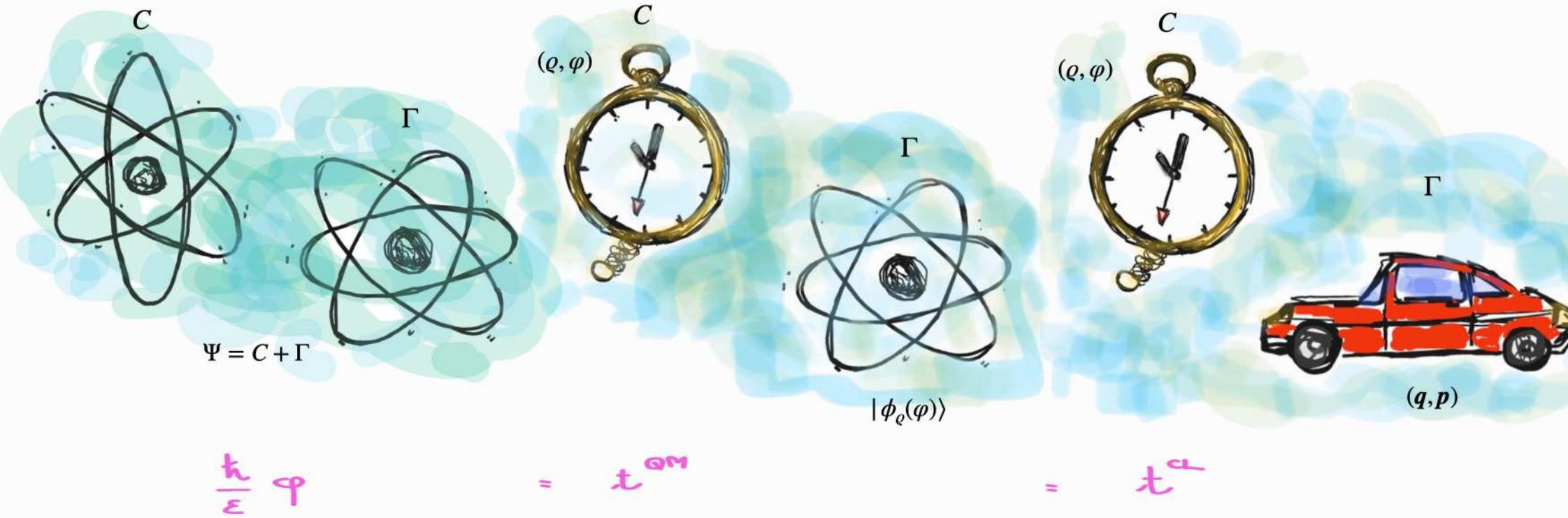
$$\{ p, H_{\Gamma} \}_r = \frac{\epsilon}{\hbar} \frac{dp}{d\varphi}$$

setting $\hbar = \hbar$

↓

!

$$t^{CL} = \frac{\hbar}{\epsilon} \varphi$$

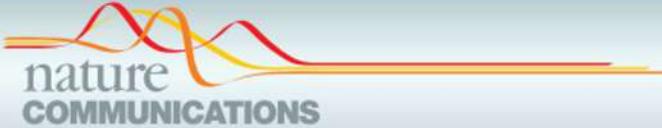


THERE IS ONLY ONE TIME

and it is a manifestation of entanglement

- energy-time uncertainty relation
- $\beta(\lambda, \omega) = \beta(e, \varphi; q, p) \rightarrow \beta(E, \varphi; q, p) \rightarrow$ spacetime ...
- Schroedinger and geodesic . . .

Refer to



ARTICLE

<https://doi.org/10.1038/s41467-021-21782-4> OPEN



Time and classical equations of motion from quantum entanglement via the Page and Wootters mechanism with generalized coherent states

Caterina Foti^{1,2,3}, Alessandro Cocco^{1,2}, Giulio Barni¹, Alessandro Cuccoli^{1,2} & Paola Verrucchi^{1,2,4}

and

PAW # 1, 4, 5

GCS # 23, 25, 40

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PARAMETRIC REPRESENTATION
WITH GCS # 24, 41